

FULL-WAVE ANALYSIS OF CASCDED JUNCTION DISCONTINUITIES OF SHIELDED COPLANAR TYPE TRANSMISSION LINE CONSIDERING THE FINITE METALLIZATION THICKNESS EFFECT

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ABSTRACT

A full-wave analysis of cascaded junction discontinuities of coplanar type transmission lines, coplanar waveguide (CPW) and finline, is implemented by mode-matching technique. Results of the frequency dependence and the transition-length dependence of the scattering parameters of the CPW-to-finline transition incorporating the finite metallization thickness effect are presented.

INTRODUCTION

For the monolithic microwave integrated circuit (MMIC) design, the coplanar type circuit is often more suitable than the double-sided circuit [1]. When the dimensions of MMIC circuit become smaller, the coupling effect between cascade discontinuities becomes significant.

The mode-matching technique has been applied to single CPW-finline discontinuities successfully [2]. In this paper, a more complex cascaded structure is analyzed by this technique. The transition-length dependence of the scattering parameters of CPW-to-finline transition is presented for the first time, with metallization thickness being considered.

MODE-MATCHING PROCEDURE

The eigenmodes of each transmission lines can be calculated by the extended spectral domain approach [3], which includes the finite metallization thickness of the coplanar type transmission line. The geometry of the problem is shown in Fig. 1. A CPW-to-finline transition consisting of cascaded discontinuities is shielded by a WR-

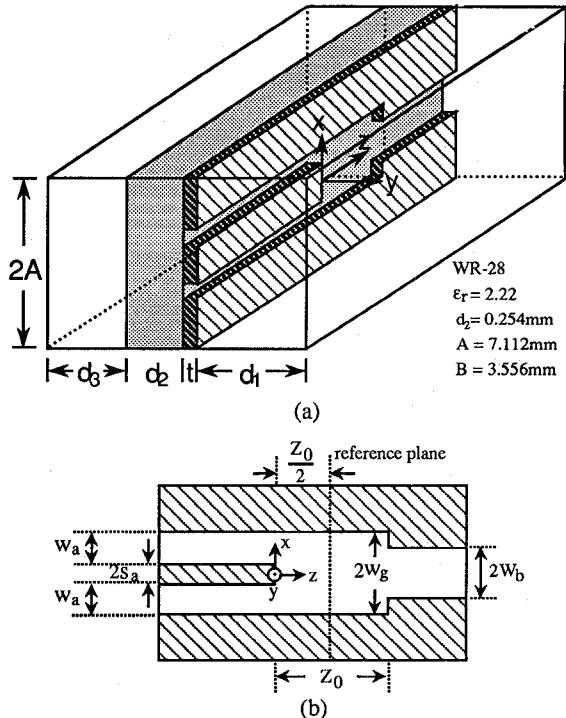


Fig. 1. (a) Shielded CPW-finline transition including the finite metallization thickness, (b) longitudinal section.

28 rectangular waveguide. The tangential field continuity at both discontinuities, $z=0$ and $z=Z_0$, can be expressed as

$$\mathbf{e}_1^a + \sum_{n=1}^{\infty} a_n \mathbf{e}_n^a = \sum_{n=1}^{\infty} c_n \mathbf{e}_n^g + \sum_{n=1}^{\infty} d_n \mathbf{e}_n^g \quad (1a)$$

$$\mathbf{h}_1^a - \sum_{n=1}^{\infty} a_n \mathbf{h}_n^a = \sum_{n=1}^{\infty} c_n \mathbf{h}_n^g - \sum_{n=1}^{\infty} d_n \mathbf{h}_n^g \quad (1b)$$

$$\sum_{n=0}^{\infty} b_n e_n^b e^{-\gamma_n^b z_0} = \sum_{n=0}^{\infty} c_n e_n^g e^{-\gamma_n^g z_0} + \sum_{n=0}^{\infty} d_n e_n^g e^{+\gamma_n^g z_0} \quad (1c)$$

$$\sum_{n=0}^{\infty} b_n h_n^b e^{-\gamma_n^b z_0} = \sum_{n=0}^{\infty} c_n h_n^g e^{-\gamma_n^g z_0} - \sum_{n=0}^{\infty} d_n h_n^g e^{+\gamma_n^g z_0} \quad (1d)$$

a, g, b : the left, center and right section of the cascaded discontinuity
 a_n, b_n, c_n, d_n : the expansion coefficient of each eigenmodes
 e_n, h_n : the normalized transverse vector electrical and magnetic fields

These simultaneous equations can be solved by carrying out power integration (3) over the cross section of the rectangular waveguide. This process results in a system of linear equations (2a-d). For the numerical computation, only a finite number of modes, N_a, N_b and N_c are used.

$$\sum_{n=1}^{N^a} a_n I_{mn}^{ag} + 0 - \sum_{n=1}^{N^b} c_n I_{mn}^{gg} - \sum_{n=1}^{N^c} d_n I_{mn}^{gg} = - I_{m1}^{ag} \quad m=1, \dots, N^a \quad (2a)$$

$$\sum_{n=1}^{N^a} a_n I_{mn}^{aa} + 0 + \sum_{n=1}^{N^b} c_n I_{mn}^{ag} - \sum_{n=1}^{N^c} d_n I_{mn}^{ag} = + I_{m1}^{aa} \quad m=1, \dots, N^a \quad (2b)$$

$$0 + \sum_{n=1}^{N^b} b_n I_{mn}^{bg} e^{-\gamma_n^b z_0} - \sum_{n=1}^{N^c} c_n I_{mn}^{gg} e^{-\gamma_n^g z_0} - \sum_{n=1}^{N^c} d_n I_{mn}^{gg} e^{+\gamma_n^g z_0} = 0 \quad m=1, \dots, N^c \quad (2c)$$

$$0 + \sum_{n=1}^{N^b} b_n I_{mn}^{bb} e^{-\gamma_n^b z_0} - \sum_{n=1}^{N^c} c_n I_{mn}^{bg} e^{-\gamma_n^g z_0} + \sum_{n=1}^{N^c} d_n I_{mn}^{bg} e^{+\gamma_n^g z_0} = 0 \quad m=1, \dots, N^b \quad (2d)$$

where

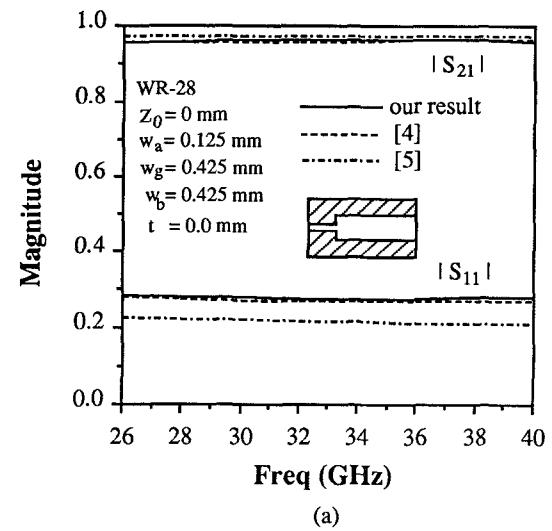
$$I_{mn}^{fg} = \int_s e_m^f \times h_n^g \cdot \hat{z} dS, \quad f, g \in \{a, g, b\} \quad (3)$$

NUMERICAL RESULTS

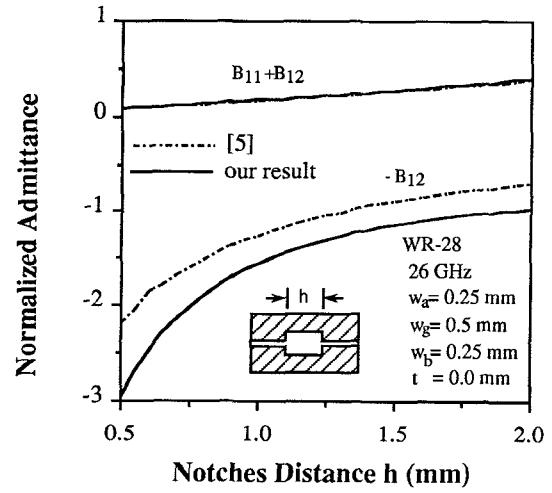
A computer code for the coplanar type transmission line cascaded junction discontinuities has been developed and checked with published results [4],[5] in Fig. 2. The necessary conditions, power conservation ($|S_{11}|^2 + |S_{21}|^2 = 1$)

and reciprocity ($S_{21} = S_{12}$), are both satisfied to within 0.001%. Six eigenmodes, i.e. $N_a = N_b = N_c = 6$, to make the results have 0.5 % convergence.

Fig 3. shows the transition-length dependence of the scattering parameters of CPW-to-fineline transition, with a reference plane at $z = Z_0 / 2$. Although the transition-length variation makes small difference in the magnitude of scattering parameters, the phase has a linear relationship with



(a)



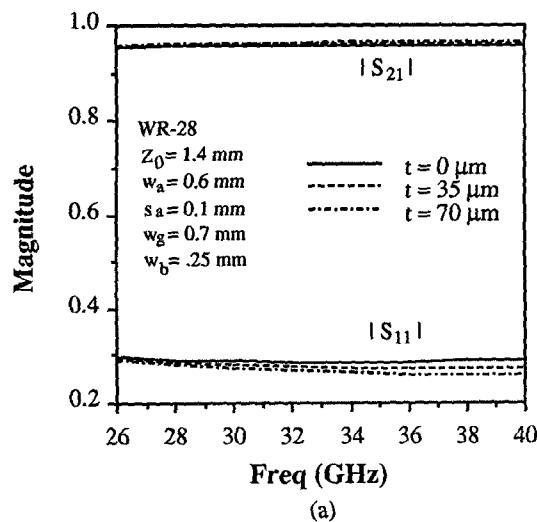
(b)

Fig. 2. Comparison of results for (a) finline step discontinuity [4],[5] and (b) inductive notches [5] with zero metallization thickness.

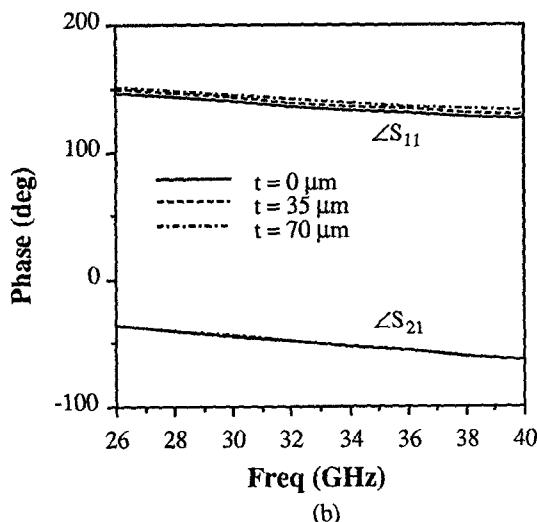
the transition-length. This phenomena is important for the impedance matching design. Fig 4. shows the frequency dependence of the scattering parameters. The linear phase variation is due to the electrical length variation at different frequencies. Fig 3 also shows an internal resonance inside the transition region and this resonance is magnified by the finite metallization thickness.

CONCLUSION

The transition-length dependence and the frequency dependence of the scattering parameters of the CPW-to-finline transition, made of two cascaded discontinuities, are presented for the first time, with metallization thickness being considered. This method can be applied to other coplanar type cascaded junction discontinuities, as CPW gap, inductive notch, capacitive strip, etc.

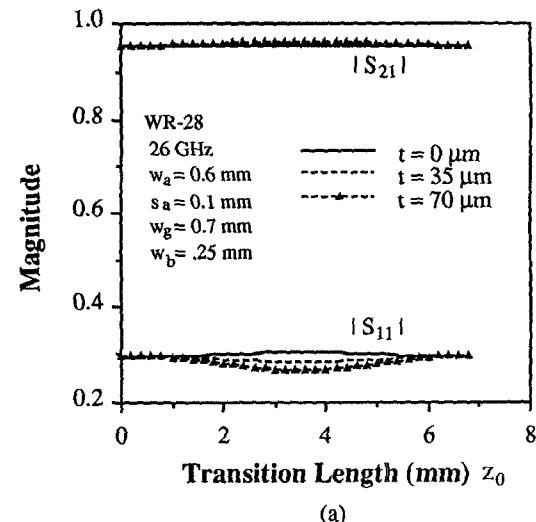


(a)

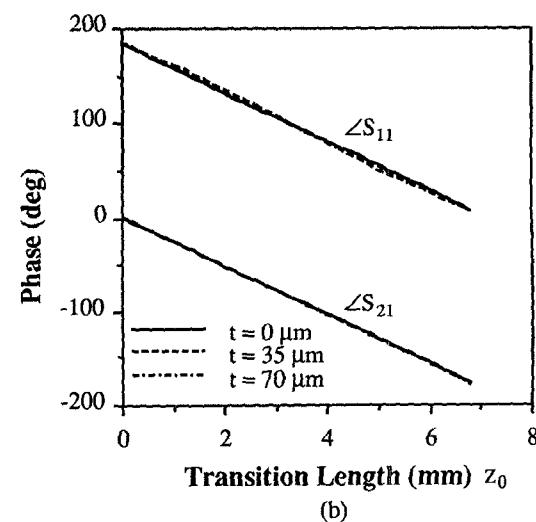


(b)

Fig. 3 Scattering parameters of a CPW-finline transition for coupled-slot mode incident at $Z_0 = 1.4 \text{ mm}$, (a)magnitude, (b)phase



(a)



(b)

Fig. 4 Scattering parameters of a CPW-finline transition for coupled-slot mode incident at 26 GHz , (a)magnitude, (b)phase

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